

 **International Research Journal of Natural and Applied Sciences ISSN: (2349-4077) Impact Factor 7.032 Volume 9, Issue 07, July 2022 Website**- [www.aarf.asia,](http://www.aarf.asia/) **Email** : [editor@aarf.asia,](mailto:editor@aarf.asia) [editoraarf@gmail.com](mailto:editoraarf@gmail.com)

**Analysis of fractional ionic character and Ionic radii for polyatomic binary compounds**

# Anil Kumar Ojha Department of Physics, Government College, Tigaon (Faridabad). E-Mail – ojha97@rediffmail.com

# **Abstract**

For investigating the nature of chemical bond, iconicity parameters and fractional ionic character parameters for polyatomic binary compounds, we have extended our knowledge from diatomic to polyatomic compounds. For calculating the values of fractional ionic character and other related parameters for polyatomic compounds the energy gap parameters  $E_g$  has been calculated. The values iconicity parameters have been obtained using the revised energy gap model due to Grimes and Cowley. Values of fractional ionic character parameters of the chemical bond are calculated for the  $A_2B$ ,  $A_3B$ ,  $AB_2$ ,  $A_3B_2$ ,  $AB_3$ and  $A_2B_3$ type polyatomic solids. In case of diatomic compounds these parameters are calculated from dielectric data .But for polyatomic compounds such data are not available. In present study we have calculated such parameters by using their universal relationship between interionic separations for the diatomic solids. The relation between fractional ionic parameters and ionic radii has also been investigated. In most cases the linear relationship between  $f_i$  and radius ratio  $r_c/r_a$  for diatomic as well as for polyatomic solids have played very useful role in understanding the relative stability for crystal structure.

**Key Words :-** Ionicity parameters ,fractional ionic character, Energy gap parameters , polyatomic compounds

# **1. Introduction:-**

The microscopic dielectric function of a solid plays the central role in the calculation of any property of the solid influenced by the screening effects of the valance and conduction electrons. Only for a small number of semiconductors, the detailed calculations have been made {1] by incorporating explicit knowledge of band structures of the particular material. In the absence of such knowledge of band structure Penn [2] derived a set of formulae based on an isotropic nearly free electron model for the energy bands. The work Penn was extended by Sriniwasan [3]. However the model adopted by Penn as well as by Sriniwasan was so complicated that the dielectric function could not be evaluated analytically. They obtained approximate expressions for the dielectric function after introducing several simplification specially the square matrix elements were replaced by interpolation formula chosen to reproduce the zero and long wave vector values and energies near the zone boundary were replaced by constant values rather than approximate expression. Grimes and Cowley [4] have rectified these shortcomings by evaluating the sum directly as integrals over the spherical Brillonin zone .Results thus obtained are sufficiently accurate at small, intermediate and large wave vectors . The modified expression for the energy gap between bonding and antibonding states obtained by Grimes and Cowley is physically more plausible and consistent with experimental data and pseudo potential calculations. Phillips and Van –

**<sup>©</sup> Association of Academic Researchers and Faculties (AARF)**

Vechten [6-7] have developed a compressive theory of iconicity of chemical bonding in binary solids using the revised energy gap expression originally due to Penn . In view of the modifications discussed above it become desirable to revise the calculations of Phillips and Van –Vechten for iconicity parameters. The studies of energy gap model and iconicity parameters have been the subject of current interest [8-10].In the present study we use the modified energy gap expression as obtained by Grimes and Cowley in order to revise the estimate of iconicity parameters due to Phillips . The revised values of iconicity parameters thus obtained for  $A^{N}B^{8-N}$  type polyatomic solids depends on radius ratio in the same manner as they found in case of diatomic solids.

# **2. Method of Analysis**

The energy gap  $E<sub>g</sub>$  can be split in to an ionic (heteropolar ) part 'C' and Covalent (homoploar )  $E_h$  such that [6]

 $E^2_{\text{g}} = E^2_{\text{h}} + C^2$ 

 (1) The homopolar part  $E_h$  has been taken to depend universally on interionic separation R as  $E_h = A R^{-k}$ 

(2)

Here we have revised values of  $A = 54.0$  J and  $k = 2.52$  given by Grimes and Cowley [4] .The constant A and k were considered by Phillips and Van-Vechten to be universally constants for different crystals which were A =63.58 and k=2.48. Values of  $E<sub>h</sub>$  for polyatomic compounds can therefore be calculated with the help of equation (1) by taking appropriate values of R . The lattice parameters or interionic separation can be divided in to two parts by introducing the concepts of ionic radii such that

 $R = r_{+} + r_{-}$ 

(3)

Where R is the nearest neighbor interionic separation and  $r_{+}$  (r.) the radii of cation (anion). In writing equation (3) it has been assumed that the ions are spherical and touching each other and this equation has been used by traditional workers [11-15] in addition to one more criteria relation to the radius ratio  $(r_{+}/r_{-})$  with the ratio of polarizabilities of effective nuclear charge of ions . The electron density measurements [12] also reveal that ions are nearly spherical and one can make use of equation 930 to obtain additive radii characteristics of ions. It is evident that ionic radii estimated by traditional workers deviate significantly from corresponding values based on electron density measurements. These deviations have been explained by Tosi [3] in terms of Fagan's rule for losing effect of cations and the tightening effect of anions in the crystalline state relative to the Free state. The ionic radii can also be related to the effective nuclear charge as

 $r_{+} = [C_n / (Z-S)_{+} ]$ 

(4)

Where  $C_n$  is a constant depending on the number of electrons in the ion.C n is same for a sequence of isoelectronic ions. S is the size screening parameter. Z is the nuclear charge (Z-S) is then known as effective nuclear charge and  $+$  and  $-$  subscripts for cation and anion respectively.

Using Pauli's relation values of  $C_n$ , the values of S can be calculated .When equation (4) is applied to a pair of isoelectronic ions yields

$$
(\mathbf{r}_{+}/\mathbf{r}_{-}) = [(Z-S) / (Z-S)_{+}]
$$
\n(5)

Equation (3) and (5) can be used to calculate radii taking experimental values of interionic separation in crystals composed of isoelectronic ions. Values of size screening parameters S have been taken to be the same as those for the corresponding isoelectronic alkali or halogen ions. Values of radii thus calculated for divalent and trivalent ions are given in table -1. The

#### **© Association of Academic Researchers and Faculties (AARF)**

radii for several other divalent and trivalent ions have also been deduced using the additivity rule (3) and experimental values of nthe lattice parameters for II-VI and III-V crystals, these values also included in Table -1.

For calculating the ionic (heteropolar) part C of the energy gap  $E<sub>g</sub>$  we have to investigate the nature of dependence of C on R in diatomic compounds for which values of C have already been determined. For doing this we first plot 'ln C' versus 'ln R' for I-VII,II-VI and III-V diatomic crystals. These plots (Figures 1-3) are linear and systematic . They are cation dependent i.e. compounds with the same cations lie on the same line. In fact we can write the following relationship

 $\widetilde{C} = BR^{-t}$ 

(6)

Where 'B' and 't' are constant for given cation .Values of B and t for various cation are given in table -2. The values of B and t thus derived have been used to predict the values of heteropolar energy 'C' for polyatomic compounds with corresponding cations. Values of heteropolar energy 'C' derived from (6) and values of homopolar energy from equation (2) are used to obtain the values of energy gap  $E<sub>g</sub>$  using equation (1) .Calculated values of 'C"  $E_g$  and  $E_h$  are reported in Table -3

The iconicity pf chemical bonds is defined in terms of fractional ionic character as given below

 $f_i = [C^2 / E^2]$ 

 $\lbrack g \rbrack$  (7) The values obtained from equations (1) and (6) can used to calcuiate the values of fractional ionic character which are reported in Table -3.

# **3. Result and Discussion**

One of the most fundamental questions in the field of solid state physics is why a given solid crystallizes in a particular structure.= In fact many important phenomenon such as magnetism, superconductivity, ferroelectricity tends to occur in a definite structure and to be absent in other structure. The splitting of energy gap  $(E_{g})$  in to ionic (hetteropolar) C and covalent (homopolar)  $E_h$  parts is made in a symmetrical manner (Equation -1) and it has found that the values of fractional ionic character parameter can provide a criteria for distinguishing between the crystal structures with different coordination numbers. Thus the fractional character parameter  $f_i$  is an important parameter related to various crystalline state properties discussed by Phillips[6]. Here we have investigated a particular useful relation between fractional ionic character f<sub>i</sub> and radius ratio of cation and anion  $r_{+}/r_{-}$  in table -4 in case of diatomic (I-VII),(II-VI) and (III-V)as well as for polyatomic binary compounds. In most of the cases the linear relationship are founds[18-21]. We have plotted some figures for polyatomic compounds. It is interesting to mention here that the parameters  $f_i$  and  $r_c$   $r_a$  have played very useful role in understanding the relative stability of the crystal structure. The values of  $f_i$  in polyatomic solids depends on the radius ratio in the same manner as they depend in case of diatomic solids [19-20]. This provides a support for the values of  $f_i$  calculated in the present study .It is the question to be investigated whether for the present study distinction could be made between the four fold and six fold coordinated structures by using the concept for fractional ionic character as it has been made for the diatomic solids [14-16].

### **Acknowledgements**

Thanks are due to Dr. B.P. Singh. Prof and Head Physics Department, Institute of Basic sciences, Khandari Agra ( Dr. B.R. University Agra campus) for his valuable guidance and discussion. Thanks are also due to Prof. Jai Shanker for his support.

**© Association of Academic Researchers and Faculties (AARF)**

**Table-1**



Values of additive radii (in  $10^{-1}$  nm) for ions

# **Table -2**

Values of constant 'B' and't' for various cations calculated from the figures [1-3]



#### **Table -3**

Calculated Values of iconicity parameters C,  $E_h$  and Energy gap Gap  $E_g$ ,  $f_i$  and radii ratio  $r_C/r_a$ 



# **© Association of Academic Researchers and Faculties (AARF)**



#### **© Association of Academic Researchers and Faculties (AARF)**



#### References

- [1] Harison W.A.; Electronic structure and properties of solids (1980)
- [2] Shanker J. and Dixit S; Phys.State sol. (a) 123.17(1991)
- [3] Phillips J.C.;Rev.of Mod.Phys.B17,42(1970 .; Phy.Rev.B29.6583(1964) ; ; Phys.State sol.(a)18,55(1966),Phys Rev.Lett.20,550(1968).
- [4] Pauling L.;Nature of Chemical Bond,CornwellUniversity,Ithaka,N.Y.(1960); ;Proc.Roy Soc.A1114,181(1927);J.Am.chem.Soc.54,3;Phjy570(1932).
- [5} P.Boguslawasi; ; Solid state.Comm.(USA),7,626(1986)
- [6] Penn D.R.; Phys.Rev. 128, 2093 (19692)
- [7] Wimples And Didomenics; Phys. Rev. Lett. 23, 1156(1969).
- [8] WimplesH.J.Chem.Phys 67,2151(1977).
- [9] Shanker J.; Goyal S.C. and Verma M.P. ; Phys. Rev. B14, 4669(1976).
- {10} Levine
- B.F.;Phy.Rev.B7,2600(1973),J.chem.Phys.,59(1973);Phys.Rev.:ett.25,44(1970).
- [11] Van. Vechten J.A. ; Phys. Rev. 89, 469(1969); Phys. Rev. B1, 3351(1970).
- [12] Grmes And Cowley F.R. ;Can.J.Phys 53,2549+ (1975).
- [13] Phillips J.C. and Van. Vechten J.A. Phys Rev. 183,709(1976).
- [14] Motfit W.F. ;Proc.Roy Soc. London A196,510(1949) ; A 202,548(1959).
- [15] Pentilides S.T.; Phys Rev. B11, 5082(1975).
- [16] Singh B.P. Phys. stat. sol. (b) 162, 329 (1990).
- [17] Shanker J.and Aharma O.P.; Philosophical Magzine 35, 6(1977).
- [18] Singh B.P.,Baghel V.S.and Bhagel K.S.;Ind.J.Pure & applied Phys.12,293(2009).
- [19] Singh B.P. and ojha A.K.;Indian J.Phys.76A(3)297(2003).
- [20] Singh B.P..Ojha A.K. And S.Tripti;Physica B,350338(2004).
- [21] Singh B.P.,Tripti S.and Singh Vipnesh Indian J.pure & AppliedPhys.46,502(2008).